

# Solving Hidden Markov Problem By Spectral Approach

C. P. Fung and B. L. Rozovskii\*  
Center for Applied Mathematical Sciences  
University of Southern California  
Los Angeles, CA 90089-1113

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**Abstract** The objective of this paper is to develop a stochastic spectral approach to solve hidden Markov problems based on a Wiener Chaos expansion. We prove that the set of Fourier coefficients in the Cameron-Martin development of the unnormalized posterior probability is a solution of a simple recursive system of infinite order deterministic equations. This decomposition separates observations and parameters which suggests a natural numerical algorithm for solving this problem.

## 1 Introduction

The objective of this paper is to develop a stochastic spectral approach to solve hidden Markov problems. Specifically we estimate the distribution of a continuous time Markov chain given noisy observations of its path. Similar approach was developed by Mikulevicius-Rozovskii [4] for nonlinear filtering of diffusion Markov processes (see also [3], [5]).

We consider the filtering scheme where the signal process  $x(t)$  is a homogeneous jump Markov process. The observation process  $y(t)$  is of the form

$$y(t) = y_0 + \int_0^t h(x(s))ds + W(t),$$

where  $W(t)$  is a Brownian motion. A fundamental objective of the filtering theory is to compute the conditional distribution of the signal process  $x(t)$  given the observations  $y(s), s \leq t$ . This problem reduces to solving an infinite system of stochastic differential equations (Rozovskii-Shiryaev [5])

$$dq_i(t) = \sum_j \lambda_{ji} q_j(t) dt + h_i q_i(t) dy(t) \quad (1.1)$$

$$q_i(0) = p_0^i \quad (1.2)$$

where  $\{p_0^i, i = 1, 2, \dots\}$  is the law of  $x(0)$  and  $h_i = h(i)$ . A solution to this equation is usually referred to as unnormalized posterior probability.

In this paper we present a spectral decomposition of a solution of (1.1), (1.2) based on the Cameron-Martin orthogonal development of  $L_2$ -functionals of a Gaussian process (Cameron-Martin [1], Hida [2]). Specifically we prove that this solution can be written in the form (see Theorem 2.1)

$$q_i(t) = \sum_{\alpha \in J} \frac{1}{\sqrt{\alpha!}} \varphi_\alpha^i(t) \xi_\alpha(y) \quad a.s. \quad (1.3)$$

where  $J$  is the set of multi-indices,  $\xi_\alpha(y)$  are products of Hermite polynomials of Wiener integrals  $\int_0^t m_j(s) dy(s)$ , where  $\{m_k\}$  is a complete orthonormal system in  $L_2(0, t)$ . We also prove that the Fourier coefficients  $\{\varphi_\alpha^i(t), i = 1, 2, \dots, \alpha \in J\}$  satisfy the recursive system of infinite order deterministic equations

$$\begin{aligned} \frac{d\varphi_\alpha^i(s)}{ds} &= \sum_j \lambda_{ji} \varphi_\alpha^j(s) + \sum_k \alpha_k m_k(s) h_i \varphi_{\alpha(k)}^i(s) \\ \varphi_\alpha^i(0) &= p_0^i 1_{\{|\alpha|=0\}}. \end{aligned}$$

It is important that decomposition (1.3) separates observations and parameters which allows the time consuming computation of deterministic Fourier coefficients  $\varphi_\alpha^i(t)$  to be performed off-line and leaves only the computationally simple stochastic part to be performed on-line. Expansion (1.3) is a double infinite sum. We study the error results from truncating the expansion (see Theorem 2.2). Many well known time discretization schemes can be derived from (1.3), for example the splitting-up approximation.

## 2 Main Results

Let  $(\Omega, \mathcal{F}, P)$  be a probability space and  $\{x(t), 0 \leq t \leq T\}$  be a homogeneous continuous time Markov process

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with transition probability density  $[\lambda_{ij}; i, j = 1, 2, \dots]$ . Let the observation process be of the form

$$y(t) = y_0 + \int_0^t h(x(s))ds + W(t) \quad (2.1)$$

where  $\sup_i |h_i| < \infty$  and  $W(t)$  is a one dimensional Brownian motion independent of  $x(t)$ . It is known that the posterior probability  $P(x(t) = i | \mathcal{F}_t^y)$  can be written as  $q_i(t) / \sum_i q_i(t)$ , where  $\{q_i(t), 0 \leq t \leq T\}$  is the unique solution of the infinite system of stochastic differential equations (Rozovskii-Shiryaev [5])

$$dq_i(t) = \sum_j \lambda_{ji} q_j(t) dt + h_i q_i(t) dy(t) \quad (2.2)$$

$$q_i(0) = p_0^i \quad (2.3)$$

where  $\{p_0^i, i = 1, 2, \dots\}$  is the law of  $x(0)$ .

Let

$$\rho(T) = \exp\left\{-\int_0^T h(x(s))dw(s) - \frac{1}{2} \int_0^T |h(x(s))|^2 ds\right\}.$$

It is a standard fact that the measure  $\tilde{P}$  defined by  $d\tilde{P} = \rho(T)dP$  is a probability measure on  $(\Omega, \mathcal{F})$  with respect to which  $y$  is a Brownian motion. Let  $J$  be the set of multi-indices. For  $\alpha \in J$ , define the following notations:  $|\alpha| = \sum_i \alpha_i$ ,  $\alpha! = \prod_k \alpha_k!$  and  $\alpha(j) = (\alpha_1, \dots, \max(\alpha_j - 1, 0), \alpha_{j+1}, \dots)$ . Let  $\{m_k(s)\}$  be a complete orthonormal system (CONS) in  $L_2(0, t)$  and  $\xi_k = \int_0^t m_k(s) dy_s$ . Further, for  $\alpha \in J$ , let

$$\xi_\alpha(y) = \prod_k \frac{H_{\alpha_k}(\xi_k)}{\sqrt{\alpha_k!}},$$

where

$$H_j(x) = (-1)^j e^{\frac{x^2}{2}} \frac{d^j}{dx^j} e^{-\frac{x^2}{2}}$$

is the  $j$ th Hermite polynomial. Then  $\{\xi_\alpha(y), \alpha \in J\}$  is a CONS in  $L_2(\Omega, \mathcal{F}_T^y, \tilde{P})$  (Cameron-Martin [1], Hida [2]). Let  $l_2 = \{(a_1, a_2, \dots) : \sum_i a_i^2 < \infty\}$ . Denote  $|\cdot|_2$  the norm in  $l_2$ .

**Theorem 2.1** If  $\{q_i(t), i = 1, 2, \dots\}$  is the solution of (2.2), (2.3), then

$$q_i(t) = \sum_{\alpha \in J} \frac{1}{\sqrt{\alpha!}} \varphi_\alpha^i(t) \xi_\alpha(y) \quad P - a.s. \quad (2.4)$$

where  $\varphi_\alpha(s) = \{\varphi_\alpha^i(s), i = 1, 2, \dots\}$  satisfies the following recursive system of deterministic equations

$$\begin{aligned} d\varphi_\alpha^i(s)/ds &= \sum_j \lambda_{ji} \varphi_{\alpha(j)}^j(s) + \sum_k \alpha_k m_k(s) h_i \varphi_{\alpha(k)}^i(s) \end{aligned} \quad (2.5)$$

$$\varphi_\alpha^i(0) = p_0^i 1_{\{|\alpha|=0\}}, \quad (2.6)$$

for  $s \leq t$ . Moreover, the series (2.4) converges in  $L_2(\Omega, \tilde{P})$ .

For  $\alpha \in J$ , denote  $d(\alpha) = \max\{k \geq 1 : \alpha_k > 0\}$  and set  $J_N^n = \{\alpha \in J : |\alpha| < N, d(\alpha) \leq n\}$ . Expansion (2.4) is a double infinite sum, denote by  $q_{N,n}(t)$  the truncated expansion which includes only terms corresponding to  $\alpha$  in  $J_N^n$ .

**Theorem 2.2** The error bound of  $q_{N,\infty}$  is given by

$$\tilde{E}|q_{N,\infty}(t) - q(t)|_2^2 \leq e^{|h|_\infty^2 t} \frac{(|h|_\infty^2 t)^{N+1}}{(N+1)!}. \quad (2.7)$$

If  $m_0(s) = \frac{1}{\sqrt{t}}$ ,  $m_k(s) = \sqrt{\frac{2}{t}} \cos(\frac{\pi k s}{t})$ , then

$$\tilde{E}|q_{N,n}(t) - q(t)|_2^2 \leq e^{|h|_\infty^2 t} \left[ \frac{(|h|_\infty^2 t)^{N+1}}{(N+1)!} + c \frac{t^3}{n} \right]. \quad (2.8)$$

For a diffusion type state process  $x(t)$ , similar result was proved in [3]. It is clear that for  $N \geq 2$  the error is of order  $O(t^3)$ . First we will demonstrate how the splitting-up approximation can be derived from (2.4) - (2.6). Let us take  $N = \infty, n = 1$  and for  $|\alpha| = k$ , write  $\varphi_k(s)$  for  $\varphi_\alpha(s)$ . Solution of (2.5), (2.6) can be approximated by

$$\varphi_k^i(s) \approx \left( \frac{s h_i}{\sqrt{t}} \right)^k p_i(s), \quad (2.9)$$

where  $p_i(s) = P(x(s) = i)$ . This can be shown by induction. If  $k = 0$ , this is the exact solution by Kolmogorov forward equation. Suppose (2.9) holds up to  $k - 1 (\geq 0)$ . From Lemma 3.1 (below), (2.5), (2.6) follows that for  $s \leq t$ ,

$$\begin{aligned} \varphi_k^i(s) &= \frac{k}{\sqrt{t}} \int_0^s \sum_j h_j \varphi_{k-1}^j(r) p_{ji}(s-r) dr \\ &\approx \frac{k}{\sqrt{t}} \int_0^s \left( \frac{r}{\sqrt{t}} \right)^{k-1} \sum_j p_j(r) p_{ji}(s-r) h_j^k dr, \end{aligned}$$

where  $p_{ij}$  is the transition probability for  $x$ . Approximate  $p_j(r) p_{ji}(s-r)$  by  $p_j(s) p_{ji}(0)$ , then the summation gives  $h_i^k p_i(s)$ , and we arrive at (2.9). When  $s = t$ ,

$$\begin{aligned} q_{\infty,1}^i(t) &\approx \sum_{k=0}^{\infty} \frac{1}{k!} \left( \sqrt{t} h_i \right)^k H_k(\xi_1) p_i(t) \\ &= \exp\{h_i y(t) - \frac{1}{2} h_i^2 t\} p_i(t), \end{aligned} \quad (2.10)$$

the last equality holds by the fact that  $\exp\{\theta z - z^2/2\} = \sum_{k=0}^{\infty} z^k H_k(\theta)/k!$ . Formula (2.10) is the splitting-up approximation.

**Corollary 2.1** The error of splitting-up method is

$$\sup_{0 \leq t \leq T} \tilde{E} |q_{\infty,1}(t) - q_t|_2^2 \leq c e^{|\lambda|_\infty^2 T} T^3.$$

### 3 Proofs

**Proof of Theorem 2.1** The theorem is proved if solution of (2.5), (2.6) exists and is unique (see [4]). This follows by induction and Lemma 3.1 below.

**Lemma 3.1** Let  $l_1 = \{(a_1, a_2, \dots) : \sum_i |a_i| < \infty\}$ . If  $u_0 \in l_1$  and  $g(t) = (g_1(t), g_2(t), \dots)$  satisfying  $\sup_{0 \leq t \leq T} \sum_i |g_i(t)| < \infty$ , then there exists a unique solution to

$$\frac{du_i(t)}{dt} = \sum_j \lambda_{ji} u_j(t) + g_i(t) \quad (3.1)$$

$$u_i(0) = u_0^i. \quad (3.2)$$

The solution satisfies  $\sup_{0 \leq t \leq T} \sum_i |u_i(t)| < \infty$  and can be written as

$$u_i(t) = \sum_k u_0^k p_{ki}(t) + \int_0^t \sum_k g_k(s) p_{ki}(t-s) ds. \quad (3.3)$$

In addition, if  $u_0 \in l_1 \cap l_2$  and  $\sup_{0 \leq t \leq T} \sum_i g_i^2(t) < \infty$ , then  $\sup_{0 \leq t \leq T} \sum_i u_i^2(t) < \infty$ .

This lemma is well understood and we omit the proof.

**Proof of Theorem 2.2** (cf. Theorem 2.2 in [3]) Since  $\{\xi_\alpha\}$  is a CONS in  $L_2(\Omega, \tilde{P})$ ,

$$\tilde{E} |q_{N,\infty}(t) - q(t)|_2^2 = \sum_{k > N} \sum_{|\alpha|=k} \frac{|\varphi_\alpha(t)|_2^2}{\alpha!}, \quad (3.4)$$

where (see [4])

$$\begin{aligned} & \sum_{|\alpha|=k} \frac{|\varphi_\alpha(t)|_2^2}{\alpha!} \\ &= \int_0^t \int_0^t \cdots \int_0^t |T_{t-s_k} h \cdots T_{s_2-s_1} h T_{s_1} p_0|_2^2 ds^k, \end{aligned} \quad (3.5)$$

where  $T_{r-s} f(i) = \sum_j f(j) p_{ji}(r-s)$  and  $ds^k = ds_1 \cdots ds_k$ . Direct computation yields

$$|T_{t-s_k} h T_{s_k-s_{k-1}} h \cdots h T_{s_1} p_0|_2^2 \leq |h|_\infty^{2k}.$$

Now we substitute the latter into (3.5), integrate and sum over  $k$  in (3.4). To prove (2.8) it is sufficient to determine  $\tilde{E} |q_{N,n}(t) - q_{N,\infty}(t)|_2^2$  which is given by

$$\sum_{l > n} \sum_{k=1}^N \sum_{|\alpha|=k, d(\alpha)=l} \frac{|\varphi_\alpha(t)|_2^2}{\alpha!}. \quad (3.6)$$

For  $|\alpha| = k$ , let  $i_1^\alpha \leq i_2^\alpha \leq \cdots \leq i_k^\alpha$  where  $\alpha_{i_j^\alpha} \neq 0$ . It is known that (see [3])

$$\varphi_\alpha^i = \int_0^t \int_0^t \cdots \int_0^t F_i(t; s^k) E_\alpha(s^k) ds^k$$

where

$$\begin{aligned} F_i(t; s^k) &= (T_{t-s_k} h \cdots T_{s_2-s_1} h T_{s_1} p_0)_i, \\ E_\alpha(s^k) &= \sum_{\sigma \in \mathcal{P}_k} m_{i_1^\alpha}(s_{\sigma(1)}) \cdots m_{i_k^\alpha}(s_{\sigma(k)}) \end{aligned}$$

and  $\mathcal{P}_k$  is the permutation group of  $\{i_1^\alpha, \dots, i_k^\alpha\}$ . Denote  $s_j^k = (s_1, \dots, s_{j-1}, s_{j+1}, \dots, s_k)$  and note that

$$E_\alpha(s^k) = \sum_{j=1}^k m_l(s_j) E_{\alpha(l)}(s_j^k).$$

Changing the order of integration we arrive at

$$\begin{aligned} & \varphi_\alpha^i \\ &= \sum_{j=1}^k \int_0^t \int_0^t \cdots \int_0^{s_{j+1}} \int_0^{s_{j-1}} \cdots \\ & \cdots \int_0^{s_2} \left( \int_{s_{j-1}}^{s_{j+1}} F_i(t; s^k) m_l(s_j) ds_j \right) E_{\alpha(l)}(s_j^k) ds_j^{k-1} \end{aligned} \quad (3.7)$$

where  $s_0 = 0, s_{k+1} = t$ . For  $m_0(s) = \frac{1}{\sqrt{t}}$  and  $m_k(s) = \sqrt{\frac{2}{t}} \cos(\frac{\pi k s}{t})$ , integration by parts gives

$$\begin{aligned} & \int_{s_{j-1}}^{s_{j+1}} F_i(t; s^k) m_l(s_j) ds_j \\ &= F_i(t; s_1, \dots, s_{j-1}, s_{j+1}, s_{j+1}, \dots, s_k) M_l(s_{j+1}) \\ & - F_i(t; s_1, \dots, s_{j-1}, s_{j-1}, s_{j+1}, \dots, s_k) M_l(s_{j-1}) \\ & - \int_{s_{j-1}}^{s_{j+1}} \frac{\partial F_i(t; s^k)}{\partial s_j} M_l(s_j) ds_j \end{aligned} \quad (3.8)$$

where  $M_l(s) = \frac{\sqrt{2t}}{\pi l} \sin \frac{\pi l s}{t}$ . The summation over  $j$  of the first two terms of (3.8) cancels out all the terms except  $-F_i(t; s_2, \dots, s_k) M_l(0) E_{\alpha(l)}(s_1^k)$  and  $F_i(t; s_1, \dots, s_{k-1}) M_l(t) E_{\alpha(l)}(s_k^k)$ , which are both zero. So, we have

$$\begin{aligned} \varphi_\alpha^i(t) &= \sum_{j=1}^k \int_0^t \int_0^t \cdots \int_0^{s_{j+1}} \int_0^{s_{j-1}} \cdots \\ & \cdots \int_0^{s_2} f_{i,l}(t; s_j^k) E_{\alpha(l)}(s_j^k) ds_j^k, \end{aligned}$$

where

$$f_{i,l}(t; s_j^k) = - \int_{s_{j-1}}^{s_{j+1}} \frac{\partial F_i(t; s^k)}{\partial s_j} M_l(s_j) ds_j. \quad (3.9)$$

Since  $\{\alpha(l) : |\alpha| = k, d(\alpha) = l\} \subset \{\alpha : |\alpha| = k-1\}$ ,

$$\sum_{|\alpha|=k, d(\alpha)=l} \frac{|\varphi_\alpha(t)|_2^2}{\alpha!} \leq \sum_{|\alpha|=k-1} \sum_i \sum_{j=1}^k \frac{k}{\alpha!} \left( \int_0^t \int_0^{s_k} \dots \int_0^{s_{j+1}} \int_0^{s_{j-1}} \dots \int_0^{s_2} f_{i,l}(t; s_j^k) E_\alpha(s_j^k) ds_j^k \right)^2.$$

Note that the summation over  $\alpha$  in the RHS of the latter inequality gives

$$\int_0^t \int_0^{s_k} \dots \int_0^{s_{j+1}} \int_0^{s_{j-1}} \dots \int_0^{s_2} f_{i,l}^2(t; s_j^k) ds_j^k,$$

(see [3]). So

$$\sum_{|\alpha|=k, d(\alpha)=l} \frac{|\varphi_\alpha(t)|_2^2}{\alpha!} \leq k \sum_{j=1}^k \int_0^t \int_0^{s_k} \dots \int_0^{s_{j+1}} \int_0^{s_{j-1}} \dots \int_0^{s_2} |f_{i,l}(t; s_j^k)|_2^2 ds_j^k. \quad (3.10)$$

Applying Holder's inequality in (3.9) we see that  $|f_{i,l}(t; s_j^k)|^2$  is bounded by

$$\frac{c}{l^2} t(s_{j+1} - s_{j-1}) \int_{s_{j-1}}^{s_{j+1}} \left( \frac{\partial F_i(t; s^k)}{\partial s_j} \right)^2 ds_j.$$

Note that  $\partial F_i(t; s^k)/\partial s_j =$

$$(T_{t-s_k} h \dots h L T_{s_j-s_{j-1}} h T_{s_{j-1}-s_{j-2}} \dots T_{s_1} p_0)_i - (T_{t-s_k} h \dots h L T_{s_{j+1}-s_j} h \dots T_{s_1} p_0)_i, \quad (3.11)$$

where  $L$  is the operator  $(Lu)_i = \sum_j \lambda_{ji} u_j$ . It is readily checked that  $L$  is a bounded linear map from  $l_1$  to  $l_1 \cap l_2$ . Let  $|Lu|_1, |Lu|_2 \leq \tilde{c}|u|_1$ . Applying the inequality  $|T_s v|_1 \leq |v|_1$  repeatedly in (3.11), we obtain

$$\sum_i \left| \frac{\partial F_i(t; s^k)}{\partial s_j} \right|^2 \leq 4\tilde{c}^2 |h|_\infty^{2k}.$$

Therefore

$$|f_{i,l}(t; s_j^k)|_2^2 \leq \frac{c}{l^2} |h|_\infty^{2k} t^2 (s_{j+1} - s_{j-1}).$$

Substituting into (3.10) and integrate, we arrive at

$$\sum_{|\alpha|=k, d(\alpha)=l} \frac{|\varphi_\alpha(t)|_2^2}{\alpha!} \leq \frac{c}{l^2} \frac{k(|h|_\infty^2 t)^{k-1}}{(k-1)!} t^3.$$

This together with (3.6) and (2.7) gives the error bound in (2.8).  $\square$

## 4 References

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